

This maple script estimates the magnitude of the elastic strain energy of a nonhydrostatically stressed isotropic solid.

The derivation corrects a math error in the expression of Connolly (Equation 29 of J.A.D. Connolly, 2009, *Geochemistry Geophysics Geosystems*, 10:10, Q10014, doi:10.1029/2009GC002540) who gives the total strain energy of an isotropic solid as

$$\begin{aligned} U = V0 * & ((e[1]^2 + e[2]^2 + e[3]^2) * c[1,1] / 2 \\ & + (e[1]*e[2] + e[1]*e[3] + e[2]*e[3]) * 2 * c[1,2] \\ & + (e[4]^2 + e[5]^2 + e[6]^2) * c[4,4] / 2) \end{aligned}$$

rather than

$$\begin{aligned} U = V0 * & ((e[1]^2 + e[2]^2 + e[3]^2) * c[1,1] / 2 \\ & + (e[1]*e[2] + e[1]*e[3] + e[2]*e[3]) * c[1,2] \\ & + (e[4]^2 + e[5]^2 + e[6]^2) * c[4,4] / 2), \end{aligned}$$

the difference being a factor of two in the second term of the second factor, where $e[i]$ and $c[i,j]$ are the elastic strain components and stiffness coefficients. The strain components $e[i]$ are related to the elements of the strain tensor $e[j,k]$ as: $e[1] = e[1,1]$, $e[2] = e[2,2]$, $e[3] = e[3,3]$, $e[4] = 2*e[2,3] = 2*e[3,2]$, $e[5] = 2*e[1,3] = 2*e[3,1]$, $e[6] = 2*e[2,1] = 2*e[1,2]$.

To make an order of magnitude assesment of the effect, non-hydrostatic (deviatoric and differential) stresses are assumed to be of magnitude $\Delta\sigma$ ("delta sigma" in Equation 30 of Connolly 2009) and the shear modulus (μ) is estimated as one third the bulk modulus (K).

Below

C - stiffness matrix

$c[i,j]$ - stiffness matrix component i,j

Epsilon1 - vector form of the strain tensor in Voigt notation.

Sigma1 - vector form of the stress tensor in Voigt notation.

DU - the total elastic strain energy (relative to the unstrained state).

strain_energy - the nonhydrostatic strain energy

the derivation follows from Callen 1985 assuming a Hookean elastic solid.

```
> restart;
> for i from 1 to 6 do;
>   for j from 1 to 6 do;
>     c[i,j] := 0 :
>   end do; end do;
> eq := 2*mu *(1 + vnu) = 3* K*(1-2*vnu);
```

```

> nu := solve (eq,vnu);
> E := 2*mu*(1+nu);

```

compliance matrix elements for an isotropic solid:

```

> c11 := 2*mu + (K-2/3*mu); c44 := mu; c12 := c11 - 2*c44;
> c[1,1] := c11; c[2,2] := c[1,1]; c[3,3] := c[1,1]; c[1,2] := c12;
c[1,3] := c[1,2]; c[2,3] := c[1,2]; c[2,1] := c[1,2]; c[3,1] :=
c[1,2]; c[3,2] := c[1,2]; c[4,4] := c44; c[5,5] := c[4,4]; c[6,6] :=
c[4,4];
> with(LinearAlgebra):
> C := Matrix(6,6,c);
> Sigma1 :=
Vector(6,[sigma[1],sigma[2],sigma[3],sigma[4],sigma[5],sigma[6]]):
> Epsilon1 :=
Vector(6,[epsilon[1],epsilon[2],epsilon[3],epsilon[4],epsilon[5],epsilon[6]]):
> Sigma := collect(simplify(C.Epsilon1),{mu,K});

```

compliance matrix S, Epsilon is the strain components in terms of stress.

```

> S := MatrixInverse(C);
> Epsilon := collect(simplify(S.Sigma1),{mu,K});

```

total strain energy/V0

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> total_strain_energy := 0:
> for k from 1 to 6 do:
> total_strain_energy := total_strain_energy +
Sigma[k]*epsilon[k]/2 :
> end do:

```

$$eq := 2 \mu (1 + vnu) = 3 K (1 - 2 vnu)$$

$$v := \frac{3 K - 2 \mu}{2 (3 K + \mu)}$$

$$E := 2 \mu \left(1 + \frac{3 K - 2 \mu}{2 (3 K + \mu)} \right)$$

$$c11 := \frac{4 \mu}{3} + K$$

$$c44 := \mu$$

$$c12 := K - \frac{2 \mu}{3}$$

$$C := \begin{bmatrix} \frac{4\mu}{3} + K & K - \frac{2\mu}{3} & K - \frac{2\mu}{3} & 0 & 0 & 0 \\ K - \frac{2\mu}{3} & \frac{4\mu}{3} + K & K - \frac{2\mu}{3} & 0 & 0 & 0 \\ K - \frac{2\mu}{3} & K - \frac{2\mu}{3} & \frac{4\mu}{3} + K & 0 & 0 & 0 \\ 0 & 0 & 0 & \mu & 0 & 0 \\ 0 & 0 & 0 & 0 & \mu & 0 \\ 0 & 0 & 0 & 0 & 0 & \mu \end{bmatrix}$$

$$\Sigma := \begin{bmatrix} K(\varepsilon_1 + \varepsilon_2 + \varepsilon_3) + \left(\frac{4}{3}\varepsilon_1 - \frac{2}{3}\varepsilon_2 - \frac{2}{3}\varepsilon_3\right)\mu \\ K(\varepsilon_1 + \varepsilon_2 + \varepsilon_3) + \left(-\frac{2}{3}\varepsilon_1 + \frac{4}{3}\varepsilon_2 - \frac{2}{3}\varepsilon_3\right)\mu \\ K(\varepsilon_1 + \varepsilon_2 + \varepsilon_3) + \left(-\frac{2}{3}\varepsilon_1 - \frac{2}{3}\varepsilon_2 + \frac{4}{3}\varepsilon_3\right)\mu \\ \mu\varepsilon_4 \\ \mu\varepsilon_5 \\ \mu\varepsilon_6 \end{bmatrix}$$

$$S := \begin{bmatrix} \frac{3K + \mu}{9\mu K} & -\frac{3K - 2\mu}{18\mu K} & -\frac{3K - 2\mu}{18\mu K} & 0 & 0 & 0 \\ -\frac{3K - 2\mu}{18\mu K} & \frac{3K + \mu}{9\mu K} & -\frac{3K - 2\mu}{18\mu K} & 0 & 0 & 0 \\ -\frac{3K - 2\mu}{18\mu K} & -\frac{3K - 2\mu}{18\mu K} & \frac{3K + \mu}{9\mu K} & 0 & 0 & 0 \\ 0 & 0 & 0 & \frac{1}{\mu} & 0 & 0 \\ 0 & 0 & 0 & 0 & \frac{1}{\mu} & 0 \\ 0 & 0 & 0 & 0 & 0 & \frac{1}{\mu} \end{bmatrix}$$

$$E := \begin{bmatrix} \frac{\frac{1}{3}\sigma_1 - \frac{1}{6}\sigma_2 - \frac{1}{6}\sigma_3}{\mu} + \frac{\frac{1}{9}\sigma_1 + \frac{1}{9}\sigma_2 + \frac{1}{9}\sigma_3}{K} \\ \frac{-\frac{1}{6}\sigma_1 + \frac{1}{3}\sigma_2 - \frac{1}{6}\sigma_3}{\mu} + \frac{\frac{1}{9}\sigma_1 + \frac{1}{9}\sigma_2 + \frac{1}{9}\sigma_3}{K} \\ \frac{-\frac{1}{6}\sigma_1 - \frac{1}{6}\sigma_2 + \frac{1}{3}\sigma_3}{\mu} + \frac{\frac{1}{9}\sigma_1 + \frac{1}{9}\sigma_2 + \frac{1}{9}\sigma_3}{K} \\ \frac{\sigma_4}{\mu} \\ \frac{\sigma_5}{\mu} \\ \frac{\sigma_6}{\mu} \end{bmatrix}$$

the bulk strain is $\text{tr}(\epsilon)/3*I$, drop I

```
> Eps_tr := (epsilon[1] + epsilon[2] + epsilon[3]) / 3;
```

eps_tr is the symbolic version of $\text{Eps_tr}/3$:

```
> dilational_strain_energy :=  
collect(subs(epsilon[1]=eps_tr,epsilon[2]=eps_tr,epsilon[3]=eps_tr  
,epsilon[4]=0,epsilon[5]=0,epsilon[6]=0,total_strain_energy),{mu,K  
});  
> deviatoric_strain_energy := collect(simplify(subs(epsilon[2] =  
-(epsilon[1]+epsilon[3]),subs(eps_tr=Eps_tr,total_strain_energy -  
dilational_strain_energy))),{mu,K});
```

$$Eps_tr := \frac{1}{3}\varepsilon_1 + \frac{1}{3}\varepsilon_2 + \frac{1}{3}\varepsilon_3$$

$$\text{dilational_strain_energy} := \frac{9K \text{eps_tr}^2}{2}$$

$$\text{deviatoric_strain_energy} := \mu \left(2\varepsilon_1^2 + 2\varepsilon_1\varepsilon_3 + 2\varepsilon_3^2 + \frac{1}{2}\varepsilon_4^2 + \frac{1}{2}\varepsilon_5^2 + \frac{1}{2}\varepsilon_6^2 \right)$$

```
> simplify(deviatoric_strain_energy);
```

$$\frac{1}{2}\mu(4\varepsilon_1^2 + 4\varepsilon_1\varepsilon_3 + 4\varepsilon_3^2 + \varepsilon_4^2 + \varepsilon_5^2 + \varepsilon_6^2)$$

```

> deviatoric_strain_energy_in_deviatoric_strain :=
collect(simplify(subs(epsilon[1]=epsilon1[1]+eps_tr,epsilon[3]=epsilon1[3]+eps_tr,epsilon[2]=-(epsilon1[1]+epsilon1[3])+eps_tr,deviatoric_strain_energy)),{mu,K});

deviatoric_strain_energy_in_deviatoric_strain :=

$$\left( 2 \varepsilon_{1_1}^2 + 2 \varepsilon_{1_1} \varepsilon_{1_3} + 2 \varepsilon_{1_3}^2 + \frac{1}{2} \varepsilon_{4_1}^2 + \frac{1}{2} \varepsilon_{5_1}^2 + \frac{1}{2} \varepsilon_{6_1}^2 \right) \mu$$


> deviatoric_strain_energy_in_stress :=
collect(subs(epsilon[1]=Epsilon[1],epsilon[2]=Epsilon[2],epsilon[3]=Epsilon[3],epsilon[4]=Epsilon[4],epsilon[5]=Epsilon[5],epsilon[6]=Epsilon[6],deviatoric_strain_energy),{mu,K});

deviatoric_strain_energy_in_stress := 
$$\begin{aligned} & \left( \frac{4}{3} \left( \frac{1}{3} \sigma_1 - \frac{1}{6} \sigma_2 - \frac{1}{6} \sigma_3 \right)^2 + \frac{2}{3} \left( -\frac{1}{6} \sigma_1 + \frac{1}{3} \sigma_2 - \frac{1}{6} \sigma_3 \right)^2 \right. \\ & - \frac{2}{3} \left( -\frac{1}{6} \sigma_1 + \frac{1}{3} \sigma_2 - \frac{1}{6} \sigma_3 \right) \left( -\frac{1}{6} \sigma_1 - \frac{1}{6} \sigma_2 + \frac{1}{3} \sigma_3 \right) + \frac{2}{3} \left( -\frac{1}{6} \sigma_1 - \frac{1}{6} \sigma_2 + \frac{1}{3} \sigma_3 \right)^2 + \frac{1}{2} \sigma_4^2 + \frac{1}{2} \sigma_5^2 \\ & \left. + \frac{1}{2} \sigma_6^2 \right) / \mu + \left( 2 \left( \frac{1}{3} \sigma_1 - \frac{1}{6} \sigma_2 - \frac{1}{6} \sigma_3 \right) \left( \frac{1}{9} \sigma_1 + \frac{1}{9} \sigma_2 + \frac{1}{9} \sigma_3 \right) \right. \\ & + \frac{2}{3} \left( -\frac{2}{9} \sigma_1 - \frac{2}{9} \sigma_2 - \frac{2}{9} \sigma_3 \right) \left( \frac{1}{3} \sigma_1 - \frac{1}{6} \sigma_2 - \frac{1}{6} \sigma_3 \right) + \frac{2}{3} \left( \frac{1}{9} \sigma_1 + \frac{1}{9} \sigma_2 + \frac{1}{9} \sigma_3 \right) \left( -\frac{1}{6} \sigma_1 + \frac{1}{3} \sigma_2 - \frac{1}{6} \sigma_3 \right) \\ & \left. + \frac{2}{3} \left( \frac{1}{9} \sigma_1 + \frac{1}{9} \sigma_2 + \frac{1}{9} \sigma_3 \right) \left( -\frac{1}{6} \sigma_1 - \frac{1}{6} \sigma_2 + \frac{1}{3} \sigma_3 \right) \right) / K \\ & + \frac{\left( \frac{4}{3} \left( \frac{1}{9} \sigma_1 + \frac{1}{9} \sigma_2 + \frac{1}{9} \sigma_3 \right)^2 + \frac{2}{3} \left( -\frac{2}{9} \sigma_1 - \frac{2}{9} \sigma_2 - \frac{2}{9} \sigma_3 \right) \left( \frac{1}{9} \sigma_1 + \frac{1}{9} \sigma_2 + \frac{1}{9} \sigma_3 \right) \right) \mu}{K^2} \end{aligned}$$


```

`sigma1[i]` are the deviatoric principle stresses, `sig_tr` is the mean stress

```

> deviatoric_strain_energy_in_deviatoric_stress := V0 *
collect(simplify(subs(sigma[1]=sigma1[1]+sig_tr,sigma[3]=sigma1[3]+sig_tr,sigma[2]=-(sigma1[1]+sigma1[3])+sig_tr,deviatoric_strain_energy_in_stress)),{dsig,mu,K});

```

to evaluate the magnitude of the deviatoric strain energy assume all independent deviatoric stresses are $\sim dsig$

```

> approximate_deviatoric_strain_energy_in_deviatoric_stress := V0 *
collect(simplify(subs(sigma[4]=dsig,sigma[5]=dsig,sigma[6]=dsig,sigma[1]=dsig+sig_tr,sigma[3]=dsig+sig_tr,sigma[2]=-(dsig+dsig)+sig_tr,deviatoric_strain_energy_in_stress)),{dsig,mu,K});

```

```

deviatoric_strain_energy_in_deviatoric_stress :=


$$\frac{V_0 \left( \frac{1}{2} \sigma_{11}^2 + \frac{1}{2} \sigma_{11} \sigma_{13} + \frac{1}{2} \sigma_{13}^2 + \frac{1}{2} \sigma_{44}^2 + \frac{1}{2} \sigma_{55}^2 + \frac{1}{2} \sigma_{66}^2 \right)}{\mu}$$


approximate_deviatoric_strain_energy_in_deviatoric_stress :=  $\frac{3 V_0 d\sigma^2}{\mu}$ 

> bulk_strain_energy_in_mean_stress := V_0 * collect(simplify(subs(sigma[1]=sigma1[1]+P,sigma[3]=sigma1[3]+P,sigma[2]=-(sigma1[1]+sigma1[3])+P,subs(epsilon[1]=Epsilon[1],epsilon[2]=Epsilon[2],epsilon[3]=Epsilon[3],subs(eps_tr=Eps_tr,dilatational _strain_energy))),{mu,K});

bulk_strain_energy_in_mean_stress :=  $\frac{V_0 P^2}{2 K}$ 

```

observing that $\mu \sim K/3$, then the non-hydrostatic strain energy is comparable to the bulk strain energy when:

```

> eq := subs(mu=K/3,approximate_deviatoric_strain_energy_in_deviatoric_stress) = bulk_strain_energy_in_mean_stress;

eq :=  $\frac{9 V_0 d\sigma^2}{K} = \frac{V_0 P^2}{2 K}$ 

```

solving this equation for $d\sigma$ gives the magnitude of the deviatoric/differential stress at which the deviatoric strain energy is comparable to the volumetric strain energy in terms of pressure.

```

> solve(eq,dsigma);


$$\frac{\sqrt{2} P}{6}, -\frac{\sqrt{2} P}{6}$$


```

as the bulk strain energy grows as P^2 the deviatoric strain energy becomes insignificant at pressures
 $> \text{sqrt}(18)*dsig \sim 4 dsig$.