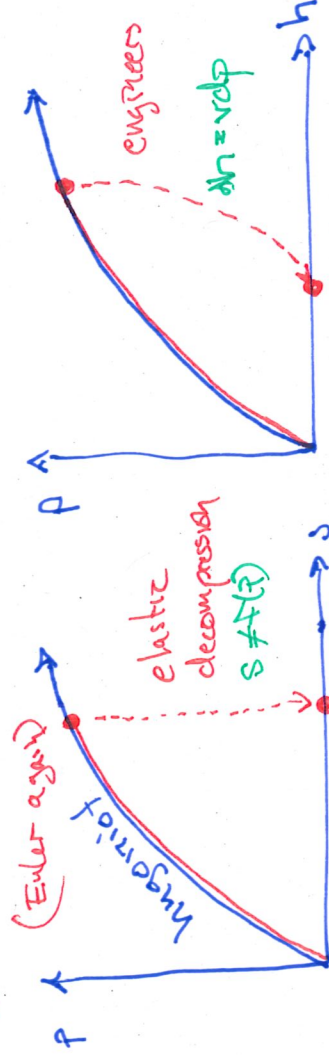


A digression: The only open adiabatic system is 2021 lecture 9.1 (Chapter 5)

one that is mechanical eq with its surrounding
 eg. mantle upwelling $h(P,S) \rightarrow h = Tds + pdv$



Two basic types of measurement:
 Mechanical Eos $V(P,T)$
 Caloric Eos $S(P,T)$

Since any phase diagram path depicts a reversible process S is always preferable to h

From theory $U(S,V,M)$ to practice (expt) $G(P,T,M) \Rightarrow$

isochemical (pure) homogeneous system $\Rightarrow G(P,T)$

$$g(P,T) = g(P,T_r) + \int_{P,T_r}^{P,T} dg$$

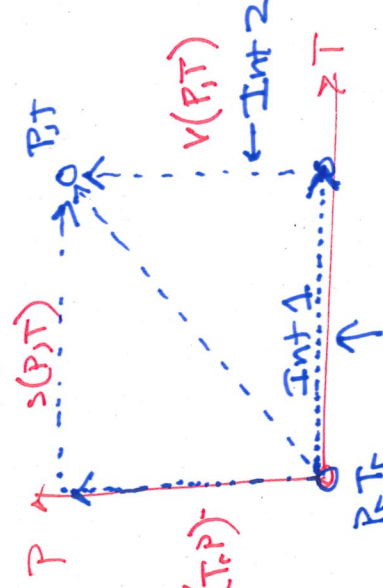
$$\int_{P,T_r}^{P,T} dg = - \int_{P,T_r}^{P,T} S(P,T) dT + \int_{P,T_r}^{P,T} V(P,T) dP$$

experiment $\Rightarrow$ caloric Eos

mechanical Eos

Caloric Eos $S(T,P_r)$

choose path to use the cheapest/best empirical data



Integral 1 The beauty is that thermodynamics provides a theoretically exact function for the unmeasured Eos

$$eg \ S(P,T) = -\partial g(P,T) / \partial T, \ c_p(P,T) \dots \quad U-(P,V)$$

Legendre Transform

Integral 1 $S(P,T)$ $\Rightarrow$

$$du = dq - dw$$

$$du = dq - pdv \Rightarrow dq = du + pdv$$

$$dh_p = Tds + vdp$$

$$\left(\frac{\partial T}{\partial P} \right)_P = \frac{1}{C_p} = \frac{1}{\left(\frac{\partial S}{\partial T} \right)_P} \Rightarrow \left(\frac{\partial T}{\partial P} \right)_P = \left(\frac{\partial S}{\partial T} \right)_P^{-1} = C_p$$

$$\Rightarrow ds = \frac{C_p}{T} dT$$

Third law entropy:

2021 lecture Q.2

$$S(P,T) = \int_0^T \frac{C_p}{T} dT \leftarrow \text{fine if no configurational entropy at } 0K$$

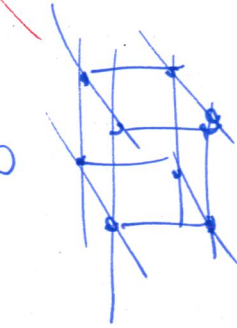
$$S(P,T) = \underline{S(P,T_r)} + \int_{T_r}^T \frac{C_p}{T} dT$$

includes
 S_{conf}

$$\begin{aligned} \text{Integral 1} &\Rightarrow \int_{T_r}^T S(P,T) dT = - \int_{T_r}^T \left(S(P,T_r) + \int_{T_r}^T \frac{C_p}{T} dT \right) dT \\ &= - \left(S(P,T_r)(T-T_r) + \int_{T_r}^T \int_{T_r}^T \frac{C_p}{T} dT \right) \end{aligned}$$

$$C_p = f(T) = a + bT + cT^2 \dots$$

Integral 2 $\rightarrow V(P,T) \rightarrow$ x-ray diffraction (a little messier for fluids)



The 1st thing we need is $V(P,T) \leftarrow$ the starting point for integral 2, ie $V(P,T)$

$$V(P,T) = V(P,T_r) + \int_{P,T_r}^{P,T} V dT \Rightarrow \text{what's provided} \Rightarrow$$

$$\alpha_P = \frac{1}{V} \left(\frac{\partial V}{\partial T} \right)_P \Rightarrow \text{expansivity (isobaric)}$$

$$= \left(\frac{d \ln V}{dT} \right)_P$$

$\Rightarrow$ simplifies notation $\begin{cases} V(P,T) = V_T \\ V(P,T_r) = V_0 \end{cases}$

$$\int_{P,T_r}^{P,T} \alpha_P dT = \int_{P,T_r}^{P,T} \frac{d \ln V}{dT} dT = \ln(V(P,T)) - \ln(V(P,T_r))$$

$$= \ln(V_T) - \ln(V_0)$$

$$= \ln(V_T/V_0)$$

$$\exp(S_{solid}) = V_T/V_0$$

volumetric strain $\rightarrow 0$

$$\text{Taylor } \exp(x) \approx 1 + x$$

$$= f_0 + \frac{df}{dx} \bigg|_{x=0} x = 1 + x$$

$$\Rightarrow \left(1 + \int_{T_r}^T \alpha dT \right) = V_T/V_0 \Rightarrow V_T = V_0 \left(1 + \int \alpha dT \right)$$

2021 lecture 9.3

V_T is the starting point for Int. 2, but the integral itself requires $V(P,T) \Rightarrow V_P \Rightarrow$ The same game

$$\beta \equiv -\frac{1}{V} \left(\frac{\partial V}{\partial P} \right)_T \Rightarrow \text{isothermal compressibility} \\ = 1/K_T \Rightarrow K_T \text{ isothermal bulk modulus}$$

$$\int_{P_r}^P \beta dP = \ln \left(V_{PT} / V_T \right) \Rightarrow V_{PT} / V_T = \exp \left(\int \beta dP \right) \\ = 1 - \int \beta dP$$

$$V_{PT} = V_T \left(1 - \int_{P_r}^P \beta dP \right) = V_0 \left(1 + \int_{T_r}^T \alpha_r dT \right) \left(1 - \int_{T_r}^{T_P} \beta dP \right) \\ = V_0 \left(1 + \int_{T_r}^T \alpha dT - \int_{T_r}^{T_P} \beta dP - \left(\int \alpha dT \int \beta dP \right) \right)$$

$$\text{Int 2} = \int_{P_r}^{P,T} V dP = \int \left(\checkmark \right) dP \rightarrow \sim 0$$

requires $\beta(P,T)$, $\alpha(P,T)$, $V_0 \Rightarrow$ gives $V(P,T)$, $\alpha(P,T)$!

In Petrology β is commonly fit to a 2nd polynomial

blows up at $P > \sim 4 \text{ GPa} \Rightarrow$ Mantle pressures require

better EOS $\Rightarrow$ Mie-Grüneisen $\Rightarrow$ isothermal EOS

Energy of a crystal lattice at cst T as a function of

volume $\Rightarrow \Omega = U - TV = A(V,T) = \sum (\text{repulsive} + \text{attractive terms})$

$T \rightarrow 0 \quad A \rightarrow U \rightarrow$ "cold" EOS (cold pressure)

$$da = -PdV - SdT \Rightarrow \left(\frac{\partial A}{\partial V} \right)_T = -P(V) \Rightarrow \text{mechanical EOS}$$

Semi-theoretical C_1/n

isothermal EOS $\rightarrow U - TS = A$



$$dA = -pdV + SdT \quad A = C_1/n - C_2/r^m$$

$$A = C_1/V^{1/3} - C_2/V^{m/3}$$

$$P = -\frac{\partial A}{\partial V} \Rightarrow P=0, V=V_0 \Rightarrow C_2 = C_1 \frac{n}{m} V_0^{(m-n)/3}$$

$$P = +\frac{C_1 n}{3V} \left(T^{n/3} - T^{m/3} \right) \Rightarrow T = \frac{V_0}{V}$$

$$K_T = -\frac{\partial P}{\partial V} \cdot V$$

$$\left. \frac{\partial P}{\partial V} \right|_{P=0} = \frac{C_1 (m-n)}{9 V_0} \Rightarrow \frac{C_1 n}{3 V_0} = -\frac{\partial K_T(0)}{m-n} \left[\frac{3}{V_0} - \frac{3}{V_0} \right]$$

$$P = -\frac{3K_T}{m-n} \left(T^{\frac{(m-3)}{3}} - T^{\frac{(n-3)}{3}} \right)$$

$$K_T = K_T(0) + K'_T P + \dots$$

$$K' = \frac{\partial K_T(0)}{\partial P} = \frac{m+n+6}{3} = 4 \leftarrow \text{experimental}$$

if $n > m$ then

$$3 < n < 9$$

$$V(V_T, K_T, K') \quad V_{T,r} = V_{T,r}^{\exp} \left(\int \alpha dT \right)$$

$$K_{T,r} = f(T) \text{ at } P_r$$

(VAP)

We need $V(P)$ but have 2 powers of V

$\Rightarrow$

Birch murnaghan

$$n=4 \quad m=2 \Rightarrow \text{to CMD}$$

Murnaghan

$$m=-3 \Rightarrow \text{to 40 GPa}$$

HP-Murnaghan

$$n=9 \quad K'=4$$

$$K = K_0 + K'P = -V \frac{dP}{dV}$$

$$\int_{P_0}^P \frac{dP}{K_0 + K'P} = \int_{V_0}^V \frac{dV}{V}$$

$$\ln \left(\frac{K_0 + K'P}{K_0 + K'P_0} \right)^{1/K'} = \frac{V_0}{V}$$

$$K_0 + K'P = \left(\frac{V}{V_0} \right)^{-K'} K_0$$

$$K'P = \left(\frac{V}{V_0} \right)^{-K'} K_0 - K_0$$

$$P = \frac{K_0}{K'} \left[\left(\frac{V}{V_0} \right)^{-K'} - 1 \right]$$

$S_0, V_0, G_0 \leftarrow$ integration constants
 $G(T), \alpha(T), K_T(T) \leftarrow$ polynomial functions of T

$\downarrow$
 $g(P, T)$

$\downarrow$
 $\alpha(P, T) = \frac{1}{V} \frac{\partial V}{\partial T} = \left(\frac{\partial}{\partial T} \left(\frac{\ln \bar{V}}{R} \right) \right)$

$$G = G_0 - \int_{T_r}^T \left(S_0 + \int_{T_r}^T \frac{G_P}{T} dT \right) dT$$

$$+ \int_{T_r, P}^{T, P} \left[\frac{V_0 \left(1 + \int_{T_r}^T \alpha(T) dT \right) \left(1 - \frac{K_P}{K_T + K'P} \right)}{\neq f(P)} dP \right]$$

$\neq f(P)$

UNITS on $\sqrt{\frac{K_{eq}}{\rho}} \Rightarrow$